# **Preferential adsorption coefficient of polymers**

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The recent formulation of the preferential adsorption coefficient,  $\lambda$ , which takes into account differences in molecular contact surface and in free volume, is tested by comparing theoretical with experimental values of  $\lambda$  from the literature. Seven different systems containing polystyrene and poly(methyl methacrylate) are considered. Agreement between theory and experiment is reached by treating the contact surface of the polymer as a fitting parameters, s. The adjusted values of s are: (a) systematically higher than the ones calculated from chain geometry; (b) largest in systems containing specific interactions (methanol). The connection between this enhanced apparent contact surface of the polymer and the ternary interaction parameter of the classical theory of  $\lambda$ , is analysed.

**Keywords** Polymer; preferential adsorption; theory; Flory-Patterson-Prigogine theory; optimization

## INTRODUCTION

In polymer-mixed solvent systems, the theoretical expression which is commonly used to interpret the experimental results of preferential sorption coefficient,  $\lambda$ , is that of Read<sup>1</sup>. This expression is derived from the Flory-Huggins (FH) thermodynamic theory of polymer solutions. According to a more recent thermodynamic theory (the so-called Flory-Patterson-Prigogine or FPP), the FH formalism is a simplification which neglects the differences in free volume and in molecular (contact) surface existing between high polymers and low molecular weight solvents. A theory of  $\lambda$  incorporating these effects of molecular surface and of free volume has been formulated recently<sup>2</sup>. It has been applied successfully to interpret preferential sorption of one of the cosolvent systems studied in our laboratory<sup>3</sup>. We extend here the application of such FPP theory of  $\lambda$  to other systems from the literature. Our aim is to test the applicability of this theory to a wide range of systems and, also, to compare it with the classical theory of Read.

Both the FH and FPP formalisms should be strictly valid for systems in which there are only weak interactions. For polar systems and for systems having specific interactions, it has been common practice to use a version of Read's formula which includes a ternary interaction parameter,  $\chi_T$ . The value of  $\chi_T$  is usually adjusted to fit the experiment. It has been suggested<sup>2,4</sup> that the effect of the ternary parameter in the classical formula is, in part, equivalent to the effects of molecular surface and of free volume in the new theory. To adjust the value of  $\chi_{\tau}$  is, therefore, a compensation for the neglect of these effects. Polar and specific interactions are additional contributions important in many systems. However, their relative importance should be appreciated fully after the effects of free volume and molecular surface are correctly taken into account by means of the new theory.

The systems which we consider here for the comparison between theory and experiment include the polymers polystyrene (PS) and poly(methyl methacrylate) (PMMA). Some of the systems show an inversion in  $\lambda$  and some are cosolvent (at room temperature). By considering a variety of different behaviours we try to draw some general conclusions regarding the best way of describing  $\lambda$ . The use of just one property is only a partial characterization of a system and an insufficient test of the theory. Therefore, we focus mainly on  $\lambda$ , but we calculate another property, the total sorption function, Y, to make a parallel comparison with  $\lambda$ .

## CALCULATION

The expression for  $\lambda$  derived from the FPP formalism is<sup>2</sup>:

$$
\lambda = -v_3 \frac{\bar{V}_1}{\bar{V}_3} \frac{l - 1 + s(\varepsilon_{13} - l\varepsilon_{23}) + (s - \alpha)\varepsilon_{12}(\varphi_1 - \varphi_2)}{l\varphi_1 + \varphi_2 - 2\varepsilon_{12}\varphi_1\varphi_2} \varphi_1 \varphi_2
$$
\n(1)

Here,  $v_3^0$  is the specific volume of polymer,  $\tilde{V}_i$  reduced volume of component  $i,\varphi$ , volume fraction of i in the solvent mixture,  $l = V_1/V_2$  is the ratio of solvent molar volumes, and s,  $\alpha$ , and  $\varepsilon_{ij}$  are parameters of the theory defined in the following way:

$$
s = S_3/S_1 \tag{2a}
$$

 $S_i$  being the molecular surface to volume ratio of component *i*. The interaction parameters  $\varepsilon_{ij}$  are:

$$
\varepsilon_{ij} = (-U_i / RT) X_{ij} / p_i^* \tag{2b}
$$

 $U_i$  being the configurational energy,  $X_{ij}$  the binary interaction parameter defined by Flory<sup>5</sup>, and  $p_i^*$  the characteristic reduction parameter for pressure. The

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Substance	$\nu$ (cm <sup>3</sup> g <sup>-1</sup> )	$\alpha$ (10 <sup>3</sup> /K <sup>-1</sup> )	$v^*$ (cm <sup>3</sup> g <sup>-1</sup> )	$p^*$ (J cm <sup>-3</sup> )	$T^*$ (K)	$S(A^{-1})$	References
$C_6H_6$	1.1462	1.223	0.8873	628	4708	0.99	4,6
CCl <sub>a</sub>	0.6304	1.229	0.4877	569	4697	0.96	4,6
$c - C_6 H_{12}$	1.2946	1.217	1.0032	531	4719	0.93	4,6
$CH_3 - C_6H_{11}$	1.3036	1.106	1.0292	450	4932	0.82	4, 6, 7
$CH3$ OH	1.2716	1.19	0.9893	500	4772	0.99	8
$\langle CH_3 \rangle_2$ CO	1.2730	1.43	0.956	627	4349	1.14	9.10
PS	0.9336	0.572	0.8098	547	7420	0.47	4,11
<b>PMMA</b>	0.808	0.575	0.701	455	7380	0.53	3

*Table I* Magnitudes characteristic of the pure substances

function  $\alpha$  is:

$$
\alpha = \alpha_1 T (p_3^* / p_1^*) (1 - T_1^* / T_3^*)
$$
 (2c)

 $\alpha_1$  being the thermal expansion coefficient of the solvent and  $T_i^*$  the characteristic reduction parameter for temperature.

For practical purposes, equation 1 can be written in terms of the usual binary interaction parameters,  $\chi_{ii}$ . According to a previous derivation<sup>3</sup>, we can take:

$$
\varepsilon_{12} \simeq \chi_{12} \tag{3a}
$$

$$
\varepsilon_{13} - l\varepsilon_{23} = \left\{ \frac{1 - l}{2} \left[ 1 - \left( \frac{\tilde{V}_3}{\tilde{V}_1} \right)^2 \right] + \left( \frac{\tilde{V}_3}{\tilde{V}_1} \right)^2 (\chi_{13} - l\chi_{23}) \right\} / \ns(s - \alpha)
$$
\n(3b)

and convert equation  $(1)$  into the more useful form<sup>3</sup>:

$$
\lambda = -v_3 \frac{H'(l-1) + H^{-1}(\chi_{13} - l\chi_{23}) + H\chi_{12}(\varphi_1 - \varphi_2)}{l\varphi_1 + \varphi_2 - 2\chi_{12}\varphi_1\varphi_2} \varphi_2 \varphi_2
$$
\n(4a)

with

$$
H \equiv (\tilde{V}_1/\tilde{V}_3)(s-\alpha) \tag{4b}
$$

$$
H' \equiv \widetilde{V}_1/\widetilde{V}_3 - [(\widetilde{V}_1/\widetilde{V}_3)^2 - 1]/2H \tag{4c}
$$

The characteristic parameters and other magnitudes needed to calculate  $V_i$ ,  $\alpha$ , and s, in equations 4 are collected in *Table 1* for all the species studied  $3,4,6-11$ . In actual calculations, for the values of  $V_1$ ,  $T_1^*, p_1^*, S_1$ , and  $\alpha_1$  to be used in equations 2 and 4, we take the mean values of the two liquids composing the solvent mixture.

The interaction parameters,  $\chi_{ij}$ , used in our calculation, have been selected from the literature  $1^{7/12-20}$  and are given in *Table 2.* For the solvent-polymer pairs we use several sets of  $\chi_{i3}$  values, denoted as A, B, C sets, in *Table 2*. For each ternary system, it is possible to decide which set best reproduces the experimental values of  $\lambda$ .

The molecular surface to volume ratio of the polymer,  $S_3$ , has a precise meaning in the theory and is obtained from the known geometry of the chain. The values for  $S_3$ corresponding to the chain geometries of PS and PMMA are given in *Table 1*. However,  $S_3$  is (usually) treated (also) as an adjustable parameter whose value is determined by fitting theory to experiment. Here we follow the same procedure and treat  $S_3$  as variable, determining its value by optimization. The criterion followed to determine  $S_3$  is

that the sum of squared deviations between experimental and calculated  $\lambda$ 's be a minimum. Since  $S_3$  appears in equation 4 always in the form of  $s - \alpha$ , it is the value of  $s - \alpha$ that is actually optimized, and from it the best value of  $S_3$ is deduced making use of equations 2a and 2c. The calculation has been programmed in a general way such that, if desired,  $S_3$  can be allowed to be a function of solvent composition. Our algorithm determines the coefficients  $a_0, a_1, a_2, \ldots$  such that the function:

$$
|s - \alpha| = a_0 + a_1 \varphi_2 + a_2 \varphi_2^2 + \dots \tag{5a}
$$

produced the least sum-of-squared deviations in  $\lambda$ . Since  $\alpha$ depends only on characteristic parameters of the pure substances, it is possible to convert the  $a_0, a_1, a_2, \ldots$ , coefficients into the expression:

$$
S_3 = C_0 + C_1 \varphi_2 + C_2 \varphi_2^2 + \dots \tag{5b}
$$

giving the optimum  $S_3(\varphi_2)$ . The values of the coefficients  $C_0, C_1, \ldots$ , obtained when we let  $S_3$  be a function of  $\varphi_2$ , are shown in *Table 2*. Also shown in *Table 2* are values of  $S_3$ determined by this same optimization method but taking  $S_3$  as constant, independent of  $\varphi_2$ . The consideration of  $S_3$ as a function of  $\varphi_2$  is an additional hypothesis, out of the theoretical considerations leading to equations 1 or 4, and should be looked on as a mere numerical fitting technique. As mentioned in the Introduction, in most of the systems the presence of polar and specific interactions renders the application of the theoretical expression, equation 4, approximate. The contributions of such interactions are, in some way, incorporated into the theory by substituting the interaction parameters  $\chi_{ij}$  appearing in equation 4 for their experimental values, as discussed previously<sup>3</sup>. But even this procedure necessitates the use of an adjustable  $S_3$  to compensate for the contributions not properly taken into account. As we shall discuss later, the difference between optimized and geometrical values of  $S<sub>3</sub>$  should be informative about the interactions in the system.

This idea of using an adjustable parameter to correct for the deficiencies of the theory, at least in what concerns its predictive capacity, is essentially the same as that which led to the extensive use of the ternary interaction parameter,  $\chi_T$ , in the classical theory. An important difference is that here we resort to this procedure only after the effects of molecular surface and free volume are explicitly taken into account, while in the classical formula of  $\lambda$ :

$$
\lambda = -v_3 \frac{\rho_0 l - 1 + \chi_{13} - l \chi_{23} + (\chi_{12} - \chi_{1}) (\varphi_1 - \varphi_2)}{l \varphi_1 + \varphi_2 - 2 \chi_{12} \varphi_1 \varphi_2} \varphi_1 \varphi_2
$$
\n(6)

Table 2 Interaction parameters used in the calculation, and results obtained for the adjustable parameters  $S_3$  and  $\chi_T$ , by least squares analysis of the experimental  $\lambda$  (equation 4 for  $S_3$  and equation 6 for  $\chi_T$ )



 $\frac{a}{2}$   $\chi_{12}$  dependent on  $\varphi_2$ ; value shown is for  $\varphi_2$  = 0.5

 $\frac{\mu}{\chi_{12}}$  = 0.303 + 0.704  $\varphi_2$  –0.057  $\varphi_2^2$ , from ref. 12

 $c$  Linearly interpolated between  $\chi_{12}$  = 1.48 ( $\varphi_1$  = 0.21) and 1.26 ( $\varphi_1$  = 0.88), from ref. 13

 $\sigma_{\chi_{12}}$  calculated from  $G^{E}$  reported in ref. 14

e From viscosity data of ref. 15

f Estimated from  $x_{12}$ ,  $x_{23}$  and the solvent composition where  $A_2$  = 0, as proposed by Dondos and Benoit<sup>16</sup>; data from refs. 16, 17

 $g$  Read<sup>1</sup>

h From solubility parameter<sup>18</sup>

<sup>*I*</sup> From the critical composition in the acetone + cyclohexanol mixture<sup>19</sup>

Calculated from data in ref. 7

 $k$  From the critical composition in the acetone + methylcyclohexane mixture<sup>13</sup>

I Average of the values in ref. 12

 $m$  Average of the values in ref. 20

 $n$  Estimate based on the poor solubility of PMMA in CCI4

 $P$  Range: 0.1842  $< \varphi_2 <$  0.9749

 $q$  Range: 0.2702  $<\varphi_2$   $<$  0.9563

everything is assigned to  $\chi_T$ .

For the sake of comparison, we also use equation 6 to compare with the experimental data of the systems listed in *Table 2*. The value of  $\chi_T$  is adjusted by optimization in the same way that is followed to determine constant  $S_3$ . The best values of  $\chi_T$  constant thus determined in each system are shown in *Table 2.* 

Our study is not limited to a single property. In addition to  $\lambda$  we also calculate  $\Delta Y$ , defined as the excess value of Y in the ternary system over its volume fraction average in the two binaries:

$$
\Delta Y = Y - \varphi_1 Y_1 - \varphi_2 Y_2 \tag{7}
$$

According to the FPP formalism,  $\Delta Y$  is given by<sup>4</sup>:

$$
\Delta Y = (\tilde{V}_1/\tilde{V}_3)^2 (s^2 - s\alpha - \alpha') \chi_{12} \varphi_1 \varphi_2 \tag{8a}
$$

where

$$
\alpha' \equiv \alpha p_3^* / p_1^* \tag{8b}
$$

However, the FH formalism obtains for  $\Delta Y$ :<sup>21</sup>

$$
\Delta Y = (\chi_{12} - 2\chi_T)\varphi_1\varphi_2 \tag{9}
$$

In each case, we calculate  $\Delta Y$  using the values of the adjusted parameters,  $S_3$  or  $\chi_T$ , deduced from the fit of  $\lambda$ data.

#### RESULTS AND DISCUSSION

The results of  $\lambda$  for the different systems are shown in *Figures 1-7.* Points correspond to experimental results<sup>1,12,13,16-18,20,22-27</sup> and curves to theoretical



*Figure 1* Preferential adsorption coefficient for the system benzene  $(1) + cyclohexane (2) + PS (3)$ . Curves  $--$ : calculated with equation 4 (0, 1,2, denoting degree of polynomial in equation 5 used to optimize  $S_3$ ); Curve  $-\cdots$ : equation 4 with fixed geometrical value of  $S_3$ . Curves  $-\cdots$ : calculated with equation 6 (0 denoting optimization with  $\chi_{\cal T}$  constant). See also text. Experimental points:  $\Box$  ref. 1 ;  $\odot$  ref. 22 ;  $\lozenge$  ref. 23 ;  $\triangle$  ,  $\vee$  ,  $\blacktriangledown$  ref. 12;  $\star$  extrapolated to infinite molecular weight in ref. 16



*Figure2* Preferential adsorption coefficient for the system cyclohexane (1) +acetone (2) + PS (3). Curves as in *Figure 1.* Experimental points extrapolated to infinite molecular weight, ref. 18

values. The continuous curves are calculated according to equation 4. The encircled numbers attached to these curves denote the degree of the polynomial in equation 5 used to optimize  $S_3$  (O is for  $S_3$  constant). The continuous curve carrying no number is for the fixed geometrical value of  $S_3$ . The broken curves are theoretical values from equation 6 calculated in two options:  $\chi_T=0$  and optimization with a constant  $\chi_T$  (broken curve 0). The curves represented correspond to the set of interaction parameters  $\chi_{i3}$  giving the best fit of equation 4 in each system. Such a set is designated as A in *Table 2.* 

# *Benzene (1) + cyclohexane (2) + PS (3): (Figure 1)*

The consideration of a variable  $S_3$  does not introduce an important improvement of the fit over the adoption of a constant  $S_3$ . A slight inversion shown by the experimental data extrapolated to  $M \rightarrow \infty^{16}$  is described by equation 4 with  $S_3$  constant but not by the classical theory with  $\chi_T$  constant. The value of  $\chi_T$  optimized with  $\lambda$ , when used in equation 9, predicts negative  $\Delta Y$ 's for  $\varphi_2$  < 0.32. However, the constant S<sub>3</sub> optimized with  $\lambda$ , when substituted in equation 8, predicts positive  $\Delta Y$ 's, for the whole range of compositions, and a maximum in  $\Delta Y$ around  $\varphi \approx 0.5$ , which is in good agreement with the experimental results of intrinsic viscosity,  $\lceil \eta \rceil^{28}$ .

#### *C yclohexane (1)+ acetone (2)+ PS (3): (Figure 2)*

The results from equation 4 are superior to the ones from the classical theory. The inversion is fairly well predicted by the optimized constant  $S_3$ . This system is cosolvent at room temperature. Its  $[\eta]$  presents a maximum at  $\varphi_2 \simeq 0.36$ .<sup>29</sup> The optimized constant S<sub>3</sub> is



*Figure3* Preferential adsorption coefficient for the system methylcyclohexane (1) + acetone (2) + PS (3). Curves as in *Figure I.*  Experimental points, ref. 13



*Figure 4* Preferential adsorption coefficient for the system ben- ' zene (1) + methanol (2) + PS (3). Curves as in *Figure I.* Experimental points:  $\Box$  ref. 25;  $\triangle$  ref. 24;  $\odot$  ref. 16 (extrapolated to infinite molecular weight)



*Figure 5* Preferential adsorption coefficient for the system carbon tetrachloride (1) + methanol (2) + PS (3). Curves as in *Figure 1.*  Experimental points extrapolated to infinite molecular weight in ref. 16



*Figure 6* Preferential adsorption coefficient for the system benzene (1) + methanol (2) + PMMA (3). Curves as in *Figure 1,* Experimental points:  $\circ$  ref. 20 ( $\lambda = v_3^0 A_1$ );  $\Box$  extrapolated to infinite molecular weight, ref. 27 (see also ref. 26)

satisfactory when used in equation 8 also, because it predicts a maximum in  $\Delta Y$  at  $\varphi_2 \approx 0.45$ . However, the optimized  $\chi_T$  which predicts negative  $\Delta Y$ 's (equation 9), is inconsistent with the  $\left[\eta\right]$  results. The use of a variable  $S_3$ improves the fit of  $\lambda$  (curve 2 in *Figure 2*) at the expense of making the predicted  $\Delta Y$  values worse, which, although still positive, present two maxima and a local minimum. As mentioned before, the use of a variable  $S_3(\varphi_2)$  is a numerical device rather than an improvement of the theory.

#### *Methyl cyclohexane (1)+ acetone (2)+ PS (3): (Figure 3)*

The curve for constant  $S_3$  is superior to the classical curve with  $\chi_T$  constant, in its description of the  $\lambda$  data. However, the opposite situation holds with respect to  $\Delta Y$ . The values of the second virial coefficient,  $A_2$ , indicate a large  $\Delta Y$  which passes through a maximum<sup>13</sup>, and this is predicted by equation 9 with the  $\chi$ <sup>T</sup> fitted to  $\lambda$ . The results of  $\Delta Y$  from equation 8 are negative due to the very low value of  $S_3$  obtained from  $\lambda$  (see *Table 2*). The fitted  $S_3$  is even lower than the geometrical value shown in *Table 1.*  This system is the only one for which the fitted  $S_3$  is smaller than its geometrical value. It is also the only for which no experimental results extrapolated to  $M \rightarrow \infty$  are available. Both facts may be correlated. Let us explore this possible correlation.

The molecular weight dependence of  $\lambda$  is of the form:  $\lambda = \lambda_{\infty} + \beta M^{-1/2}$ . The data in *Figure 3* are for <sup>13</sup>  $M_w = 1.56 \times 10^5$ . In the closely related system cyclohexane + acetone + PS, the term  $\beta M^{-1/2}$  at  $\varphi_2 \simeq 0.6$ (where  $\lambda$  is highest in *Figure 3*), corresponding to  $M = 1.56 \times 10^5$ , is larger than  $\lambda_{\infty}$ . This means that if the present system (with methyl cyclohexane) shows a similar M dependence, on extrapolation to  $M \rightarrow \infty$ ,  $\lambda$  could be reduced to less than half the value it shows in *Figure 3.* The use of non-extrapolated, large  $\lambda$  values, may be the cause of the anomalously low  $S_3$  fitted value. The value of  $\chi_T$  is also peculiar in that it is the only one found to be negative.

#### *Benzene (1) + methanol (2) + PS (3): (Figure 4)*

Only the data extrapolated to  $M \rightarrow \infty^{16}$  have been used in the fit. Here fitting a constant  $S_3$  or a constant  $\chi_T$  leads to similar curves which describe the data fairly well. However, only the optimized constant  $S_3$  gives satisfactory results for  $\Delta Y$ . The  $\Delta Y$ 's from  $S_3$  are positive and with a maximum at  $\varphi_2 \approx 0.14$ , while the  $\Delta Y$ 's from  $\chi_T$ are negative for  $\varphi_2 > 0.09$ . The use of quadratic  $S_3$  leads to a perfect fit of  $\lambda$  and positive  $\Delta Y$ 's.

#### *Carbon tetrachloride (1) + methanol (2) + PS (3): (Figure 5)*

The results are very similar to the ones discussed in the preceding system.  $S_3$  constant and  $\chi_T$  constant produce equivalent fits of  $\lambda$ .  $S_3$  quadratic gives a perfect fit. The comparison with  $\Delta Y$  is also analogous to the previous system.  $S_3$  constant leads to a maximum in  $\Delta Y$  at  $\varphi_2 \simeq 0.08$  while  $\chi_T$  leads to negative  $\Delta Y$ 's. The quadratic  $S_3$  leads also to  $\Delta Y < 0$ . Therefore,  $S_3$  constant is the best simultaneous description of  $\lambda$  and  $\Delta Y$ .



*Figure 7*  Preferential adsorption coefficient for the system carbon **tetrachloride** (1) + methanol (2) + PMMA (3). Curves as in *Figure I.*  Experimental points extrapolated to infinite molecular weight, ref. 17

 $C_6H_6 + CH_3OH$  +  $CH_3OH$   $\text{ICCl}_4 + \text{CH}_3\text{OH}$   $\text{PMMA}$   $\text{---}$  +  $\text{---}$ 



*Table 3* Description of preferential and total sorption using a constant adjustable parameter:  $S_3$  in the FPP formalism and  $\chi_{T}$  in the classical FH one. (Plus and minus signs indicate fair and poor description of the experimental results, respectively.)

## *Benzene (1)+ methanol (2)+ PMMA (3): (Figure 6)*

The description of  $\lambda$  by means of a constant  $S_3$  is poor and quadratic  $S_3$  is needed for a satisfactory fit.  $\chi_T$  gives a very good fit, almost as good as the one with quadratic  $S_3$ . In contrast to the systems considered thus far, it is not possible to find a simultaneously good description of  $\lambda$ and  $\Delta Y$  with a single constant parameter, be it  $S_3$  or  $\chi_T$ . The description of  $\lambda$  is much better with  $\chi$ <sub>T</sub> constant than with  $S_3$  constant, but the reverse is true with regard to  $\Delta Y$ . Constant  $\chi_T$  gives  $\Delta Y < 0$ , which contradicts experiments showing that  $[\eta]$  passes through a maximum at  $\varphi \simeq 0.18^{20}$ . The predictions with  $S_3$  constant are, however, correct (at least qualitatively) because they give  $\Delta Y > 0$ with a maximum at  $\varphi_2 \approx 0.33$ . The quadratic  $S_3$  which gives a good fit of  $\lambda$  predicts a maximum in  $\Delta Y$  at  $\varphi_2 \approx 0.18$ , but it is a local one,  $\Delta Y$  starting to rise again at  $\varphi_2 \simeq 0.3.$ 

#### *Carbon tetrachloride (1) + methanol (2) + PMMA (3): (figure 7)*

The results are analogous to those in the preceding system. Constant  $\chi_T$  fits well with  $\lambda$  but produces negative  $\Delta Y$ . Constant S<sub>3</sub> fits poorly with  $\lambda$  but predicts  $\Delta Y > 0$ with a maximum at  $\varphi_2 \simeq 0.35$ , agreeing (at least qualitatively) with the experimental  $[\eta]$ , showing a maximum at  $\varphi_2 \simeq 0.2^{17}$ . S<sub>3</sub> quadratic fits very well and predicts a maximum in  $\Delta Y$  at  $\varphi_2 \simeq 0.23$ , but it is a local one.

The summary of the systems studied is shown in *Table 3,*  where the plus and minus signs denote fair and poor descriptions of the data, frespectively. As we can see, the PS systems are better described with the FPP formulas than with the classical ones. The two systems containing PMMA are not well described by either theory. However, previous analysis of the system acetonitrile  $(1)$ +chlorobutane (2)+PMMA (3) showed that  $S_3$ constant gives good agreement both with  $\lambda$  and  $\Delta Y$ , while  $\chi_T$  does not<sup>3</sup>. The simultaneous presence of PMMA and methanol renders the two systems more difficult to describe.

This type of analysis depends critically on the interaction parameters chosen and, to a certain degree, on the fitting procedure. In our work,  $\lambda$  is fitted directly;  $\Delta Y$ being calculated using the parameters obtained from  $\lambda$ . The direct fit of  $\lambda$  follows a least-squares criterion. In some work, the adjustable  $\chi_T$  has been determined by solving  $\lambda = 0$  at the inversion point. This is simple, but overestimates the weight of such an inversion point. The least-squares criterion does not rely on any single point but considers all of them on the same basis.

The comparison established in *Table 3* between the FH and FPP formulas is for constant  $S_3$  and  $\chi_T$  parameters. It is sometimes found that a composition dependent  $\chi_T$  is needed to describe experimental data.  $\chi_{12}$  also varies with  $\varphi_2$  in most systems. A formula for  $\lambda$  taking into account, explicitly, such  $\varphi_2$  dependence has been developed in the FH scheme<sup>21</sup>. It should be a better description of  $\lambda$  to develop an equivalent formula in the FPP scheme by incorporating, explicitly, the derivatives of  $\chi_{12}$  with respect to  $\varphi_2$ . Our present analysis is more crude, simply using equations 4 and 6 with  $\chi_{12}(\varphi_2)$ . This avoids the calculation of  $d\chi_{12}/d\varphi_2$  and  $d^2\chi_{12}/d\varphi_2^2$ , which is impractical in many systems due to the uncertainty in the interaction parameter values.

## CONTACT SURFACE AND TERNARY INTERACTIONS

The values of constant  $S_3$  in *Table 2* are systematically larger than the geometrical values in *Table 1* (with the exception of methyl cyclohexane + acetone + PS, whose anomalous origin has been discussed previously. The value of  $S_3$  obtained previously<sup>3</sup> for acetonitrile + chlorobutane + PMMA  $(S_3 = 0.63 \text{ Å}^{-1})$  is also higher than the geometrical one. The systems containing methanol have  $S_3$  values in the range 0.73-0.84\* which are higher than the values for the other systems, which have the range 0.62-0.65\*. It seems, therefore, that for the systems having stronger specific interactions (alcohol), there is correspondingly a higher  $S_3$ . We are led to the conclusion that the strong interactions, which are not adequately taken into account by the thermodynamic theory, appear as a larger value of contact surface of the polymer. The effect of such interactions is, then, to produce an apparently larger exposure of the polymer molecular surface to solvents.

Consider, as a working hypothesis, the idea that strong interactions contribute to a larger  $S_3$ , and let us see what consequence this has on the solvation of the polymer chain. We write

$$
s = s_g + \Delta s \tag{10}
$$

with  $s_a$  meaning the geometrical value of s and  $\Delta s$  the difference between  $s<sub>g</sub>$  and the empirically adjusted s. We substitute s given by equation 10 into the expression for total sorption (equation 8), and get:

$$
\Delta Y = (s_g^2 - s_g \alpha - \alpha') (\tilde{V}_1/\tilde{V}_3)^2 \chi_{12} \varphi_1 \varphi_2 (1 + \Delta s \cdot t)
$$
\n(11a)



with

$$
t \equiv (2s_g - \alpha + \Delta s)/(s_g^2 - s_g \alpha - \alpha') \tag{11b}
$$

In equation 11,  $\Delta Y$  is the result of two contributions. The first due solely to weak interactions  $(s = s_a)$ , and the second (the one in  $\Delta s$ ) is an additional contribution arising from the enhanced apparent molecular surface of the polymer, presumably due to stronger interactions. We can see that this additional contribution is positive and increases solvation of the polymer  $(2s_q - \alpha + \Delta s)$  is always found to be positive, of the order of unity). The actual total sorption is larger than that predicted for the same system having only weak, random interactions. This is an acceptable result.

Going back to *Table 2* we analyse what happens with the classical theory. We see that  $\chi_T$  is larger in the systems containing methanol. The need of a non-zero  $\chi_T$  has usually been interpreted as being due to the influence of ternary contacts, the non-random nature of the interactions, the gathering of solvent molecules around the polymer chain, etc. Intuitively, one would think these effects would be favourable to polymer solvation. However, the role played by  $\chi_T$  in equation 8 is just the opposite to this intuitive expectation.  $\chi_T$  appears in  $\chi_{12} - \chi_T$  as a contribution opposing solvation. By writing  $\chi_T = (\gamma/2)\chi_{12}$ , we get from equation (8):

$$
\Delta Y = \chi_{12}\varphi_1\varphi_2(1-\gamma) \tag{12}
$$

which is of the same form as equation 11 but with the contribution from the ternary term paradoxically reducing total sorption.

We can now see that this paradoxical result happens because  $\chi_T$  includes not only the effects of strong interactions but also the effects of contact surface and free volume. These two latter effects make  $\Delta Y$  much smaller than predicted by FH theory  $(s_a^2 - s_a - \alpha'$  in equation 11 attains small values of the order of 0.1 or less). Once these are taken into account separately, the strong interactions can appear as a positive contribution to sorption ( $\Delta s > 0$ ) in equation 11).

There are systems with favourable 1-2 interactions contributing negatively to  $\Delta Y$ , but none of the systems considered here fall in such a category. All of the liquid mixtures have  $G^E > 0$  and solvation of the polymer coil is enhanced in the mixture.

The action of  $\Delta s$  on  $\lambda$  is more difficult to analyse because  $\lambda$  is a complex function of  $s = s_a + \Delta s$  in equation 4. If we look solely at the term  $(s - \alpha) \chi_{12}(\varphi_1 - \varphi_2)$ , we see that  $\Delta s$  reinforces the adsorption of 1 for  $\varphi_2 > 0.5$  and the adsorption of 2 for  $\varphi_2$ <0.5. That is,  $\Delta s$  favours the adsorption of the minor component in the liquid mixture thus giving a better solvent inside the coil, which is a reasonable result. However, the action of  $\chi_T$  in the term  $(\chi_{12} - \chi_7)(\varphi_1 - \varphi_2)$  of equation 6 is the opposite, because  $\chi_T$  favours the adsorption of the major component.

Our discussion of As has been purely *ad hoc.* To take into account strong interactions in a proper way would require devising a model for such interactions, in the FPP formalism, and rederiving the equations for  $\lambda$  and  $\Delta Y$ .

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